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# Reviewer Comments:

# Trends in field Quality Along the Production of the LHC dipoles and differences among Manufacturers

E. Todesco, B. Bellesia, P. Hagen, C. Vollinger

WERE

Abstract—More than two thirds of the dipoles of the Large Hadron Collider have been manufactured and their magnetic field has been measured at room temperature. In this paper we make a review of the trends that have been observed during the production. In some cases, the trends have been traced back to displacements of conductors with respect to the nominal lay-out. The analysis allows detecting the most critical zones in the superconducting coil as far as field quality is concerned. The second part of the paper makes the point of the observed differences in field quality between the three manufacturers. The analysis allows evaluating which multipoles are more affected, what magnitudes of displacements are necessary to explain these differences (the manufacturers all producing the same baseline), and what could be the origin of such differences.

HEEDED Index Terms—LHC, Superconductivity, Field Quality.

I. INTRODUCTION

o guide the particle beams in the Large Hadron Collider O guide the particle beams in the Large American (LHC) [1] 1232 superconducting dipoles with a nominal field of 8.3 T are being produced. More than ¾ of the superconducting cails have already been manufactured. One of the requirements imposed by the beam dynamics is that the deviations of the magnetic field from ideal are of the order of 10-100 ppm [2]. In this type of magnets [3],[4], the shape of the magnetic field is mainly determined by the position of the conductors, which have to satisfy tolerances better than 0.1 mm, and by the superconducting properties of the cable. As in the case of the previous superconducting magnet mass production (the Relativistic Heavy Ion Collider [5]) magnetic measurements are carried out at room temperature over 100% of the production, and over a sample of 15-20% in operational conditions at 1.9 K. Link the magnetic field quality is steered towards the beam dynamics targets [6]. Magnetic measurements are also a powerful tool to carry out a quality control since anomalies in the field can be traced back to assembly errors or faulty components [7],[8].

In this paper we review the results of magnetic measurements at room temperature of 844 dipoles of the LHC, relating the magnetic field properties to the assembly procedures and components [9-12]. The aim of the work is to

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show what degree of homogeneit has been obtained in this mass production, what are the most important mechanisms driving the magnetic field at room temperature, and the relation with the design tolerances. We recall that in the LHC dipole, field harmonics are by far dominated by the geometric contribution, since the persistent current and saturation are - ? EXPLAIN BETTER: THEIR EXCECT well reproducible

II. DIPOLE COMPONENTS, MANUFACTURERS AND ASSEMBLERS

The main dipoles of the LHC are made up of superconducting coils clamped in austenitic steel collars inside an iron yoke, the whole structure being contained in a shrinking cylinder (see Fig. 1). The cables are 15.1 mm wids, NbTi cooled to 1.9 K,. Two layers with different type of cables are would to form the coil pole. Four copper wedges are used to approximate a cos  $\theta$  geometry, thus dividing the coil in six conductor blocks (see Fig. 2). Compared to previous projects, this is the first dipole design with two apertures in a common collar (twin design), and therefore the issue of possible coupling of the magnetic fields between the two apertures is of relevance. The main component suppliers, and assemblers of the dipole are the following enes

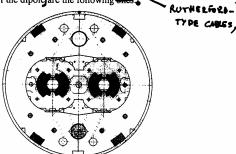


Fig. 1: Cross-section of the main LHC dipole

- Superconducting cables are manufactured by 2 different suppliers (labeled B,E) for the inner layer and 5 (coded B.C.G.D.K) for the outer one [13].
- Copper wedge spacers are manufactured by one company
- Collars are manufactured by two suppliers S1 (5/8 of the production) and S2 (3/8) [11], [12].

Iron laminations are manufactured by two suppliers (5/8 by S2 and 3/8 by S3, for the straight part).

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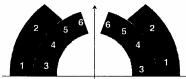


Fig. 2: Cross-section of the LHC dipole coil (upper half of one aperture) and

- The assembly of the magnet, involving the coil winding and curing, the collaring, the assembly of the yoke, and the welding of the cylinder is carried out by three manufacturers (dipole assemblers) denoted by Firm1, 2 and 3. The assembly of the collared coil follows different procedures in each assembler [11], [12]:
- o In Firm1 the pre-assembly is done with the collared coil in a vertical position, whereas in THEY ARE Firm2 and Firm3 it is horizontal.

In Firm1, collars are mounted using two 0 configurations (up down flip); in Firm3 they are mounted in two other configurations (rotation by 180 degrees); in Firm2 the four configurations are used. More details are given in [11].

Systematic differences in the coil geometry can be due to

- the cable manufacturer, in case of systematic differences in the cable size and mechanical properties [10];
- the collar manufacturer, in case of systematic differences in the collar geometry [11];
- the cold mass assembler, in case of systematic differences in the tooling used for the coil curing, in the procedure for the assembly of the collars and for the collaring [11].

Another difference in the dipoles stems from the tuning of the lay-out that North Sen carried out during the production to better match the beam dynamics targets [6]. Three crosssections have been produced: the baseline (cross-section 1, also denoted by version V6-1 [14]) Nas-Sen replaced by a new version with a different geometry of the copper wedges of the inner layer (cross-section 2 [6]), and then a mid-plane shim of 0.125 mm thickness Missen added (cross-section 3 [6]). 34 magnets have cross-section 1, 145 cross-section 2, and the rest of the production is with cross-section 3.

# III. MAGNETIC MEASUREMENTS

The magnetic field in a dipole is expressed as a power series

$$B_{y}(x, y) + iB_{x}(x, y) = 10^{-4} B_{1} \sum_{n=1}^{\infty} (b_{n} + ia_{n}) \left(\frac{x + iy}{R}\right)^{n-1}$$

where (x,y) are the transverse coordinates, R is the reference radius (17 mm in our case), and B, is the main dipolar component. The harmonics terms  $b_3, b_5, b_7 \dots$ , are generated by a coil lay-out that satisfies the up-down and left-right symmetry ("allowed" components), whereas the other harmonic terms are due to a break of these symmetries ("not allowed" components). The field harmonics are expressed in units, and one has  $b_i=10^4$  by definition. The main component

and the field harmonics are measured at room temperature ith a rotating coil of 750 mm length. 20 consecutive positions are measured, along the 14.3 m dipole. Position 1 and 20 cover the ends of the coils, and 2 to 19 the so called straight part. In this paper we will neglect the variations of the field harmonics along the magnet axis, and each dipole will be characterized by two values of the average harmonics (usually called integral), one for each aperture. These values are averages weighted with the main field component.

Measurements are carried out at the manufacturer [15] at two different stages of the assembly procedure, namely after the collaring (collared coil, i.e. the superconducting coils clamped in the collars), and after the welding of the shrinking cylinder (the so called cold mass, i.e. the collared coil inside the iron yoke and the stainless steel cylinder). Here we will present the results relative to the cold masses. Indeed, the magnetic effect of the iron yoke is very reproducible and therefore the collared coil measurement already contains the bulk of the information about the coil geometry.

Measurements are carried out in operational conditions (1.9 K) at CERN on 15% of the production, the sampling having been 100% in the pre-series phase to build solid warm-cold correlations [16]. In the mature phase of the production the sampling is around 10%. A summary of the available measurements at room temperature and at 1.9 K is given in Table I. This paper focuses the features that can be deduced by room temperature measurements.

TABLE I: NUMBER OF MAGNETS MEASURED AT ROOM TEMPERATURE AND AT

1.9 K FOR EACH DIPOLE ASSEMBLER AND CROSS-SECTION TYPE										
	Ro	om ter	nperat	ure	1.9 K					
Firm	Xs1	Xs2	Xs3	all	Xs1	Xs2	Xs3	all		
1	12	58	174	244	13	26	14	53		
2	10	31	182	223	9	13	17	39		
3	9	51	317	377	8	18	. 37	63		
all	31	140	673	844	30	57	68	155		

ARE YOU SURE? IV. BENDING STRENGTH

A. Transfer function The transfer function is defined as the average main field in the straight part of the dipole (position 2 to 1) divided by the current. The moving average per cold mass assembler has been stable after the last change of cross-section (see Fig. 3), and is oscillating along the production in a range of ±15 units. We WHAT'S remind the reader that 15 units of main field can be generated THE by a deviation in the coil diameter of 0.04 mm, which is of the same order of magnitude as the tolerances on the collar radius (0.03 mm). Since the beginning of the production, we have a systematic difference of around 10 units between Firm3 and Firm1, Firm2 being in between.

## B. Magnetic length

The magnetic length is defined as the integral of the transfer function divided by the transfer function in the straight part. NTIE moving average per manufacturer has been stable within 1-2 units since the beginning of the production (see Fig. 4). One observe a small systematic difference of 10 units, corresponding to 14 mm, between the dipoles assembled in

ALSO BEFORE!

Firm3 and in Firm1, Firm2 being in between Part of this difference is due to the iron yoke in the collared coils the difference in magnetic length between Firm3 and Firm1 is smaller (5 units).

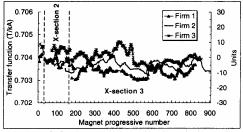


Fig. 3: Transfer function measured at r.t. along the production: moving average for each cold mass assembler.

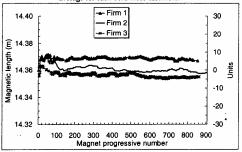


Fig. 4: Magnetic length measured at r.t. along the production: moving average for each cold mass assembler.

#### C. Integrated transfer function

The integrated transfer function is the product of the magnetic length times the transfer function. This quantity gives the actual bending strength of the dipoles, and is the only one relevant for the beam. In the technical specification it was foreseen to have a fine tuning of the length of the iron laminations to reduce the spread of the bending strength between the assemblers. This method has been tested successfully in three dipoles at an early stage of the production. Indeed, the systematic differences between Firm1 and Firm3 in the transfer function and magnetic length compensate each other, and the difference of the bending strength average between the assemblers is below 3 units. This accidental compensation has avoided the necessity of carrying out the fine tuning with the iron laminations. To date, the standard deviation of the bending strength of all manufactured dipoles at room temperature is 5.5 units, well below the specification of 8 units [2]. This corresponds to a spread in the coil radius over all the production of 15  $\mu m$  (one standard deviation). The initial installation baseline was to allocate the same dipole assembler in the same octant, to minimize the spread: the absence of significant systematic differences between assemblers has allowed to relaxing this constraint.

#### V. ALLOWED MULTIPOLES

Average and standard deviation of the odd normal multipoles are given in Table II, where we analyzed only the dipoles with cross-section 3 to have a homogeneous sample. The motivations for the cross-section changes and their effectiveness have been already presented in [6]. Here we focus our analysis on the systematic differences between dipole assemblers, which are negligible for the  $b_3$ , and rather large for  $b_5$ ,  $b_7$  and  $b_9$  (see Table II, last two lines). The normal decapole  $b_5$  in Firm1 is 0.75 units larger than in Firm2-3, and  $b_7$  in Firm2 is 0.3 units smaller than in Firm1-3. The mechanism that drives the difference in  $b_5$  is probably also affecting  $b_9$  which is 0.09 units higher in Firm1.

TABLE II: AVERAGE AND STANDARD DEVIATION OF ODD NORMAL MULTIPOLES

WILA.	b3			b5		7	b9	
Firm	ш	-	ш	σ	ш	σ	u	•
1	3.0	0.9	0.24	0.27	1.02	0.07	0.50	0.02
2	2.1	1.2	-0.54	0.38	0.73	0.12	0.41	0.02
3	3.1	0.8	-0.44	0.18	0.99	0.06	0.41	0.02
all	2.7	1.0	-0.25	0.42	0.91	0.15	0.44	0.04
1-3	-0.1		0.68		0.03		0.09	
2-3	-1.0		-0.10		-0.26		0.00	

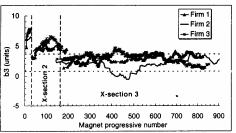


Fig. 5: Multipole b, measured at r.t. along the production: moving average for each dipole assembler (markers) and targets for systematic (dotted line).

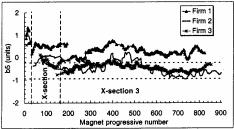


Fig. 6: Multipole b, measured at r.t. along the production: moving average for each dipole assembler (markers) and targets for systematic (dotted line).

These differences are also visible in the previous crosssections (see Figs. 5-7). They cannot come from the collar supplier, since Firm1 and Firm2 use the same supplier S1. The best candidate for the difference between Firm2 and Firm3 is a 0.075 mm inward radial shift of block5 (see Fig. 8) in all quadrants, which gives a large effect on  $b_7$  (-0.27 units), without affecting  $b_5$ , with a relatively small impact on  $b_3$  (-1.6 units), close to what has been measured. On the other hand, no simple movement of a conductor block can account for the systematic differences between Firm1 and Firm3. Indeed, the order of magnitude of displacements that can give the measured difference of 0.7 units of  $b_5$  (see Table II) is also in this case of 0.05 to 0.1 mm, i.e. not far from the tolerances of the design.

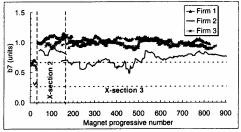


Fig. 7: Multipole b, measured at r.t. along the production: moving average for each dipole assembler (markers) and targets for systematic (dotted line).

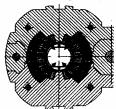


Fig. 8: Coil displacement of block5 that could account for the difference in allowed multipoles measured at r.t. between Firm2 and Firm3 (not in scale).

Using the approach developed in [17], we find that the measured spread of the allowed multipoles in each assembler matches the spread due to a random movement of the conductor blocks of 0.05 mm amplitude (one sigma). Neither the copper wedges [9] nor the collars [11] tolerances are found to be the driving mechanism of the allowed multipoles. We finally point out that a fairly good correlation (r=0.7-0.8) is found between the apertures [11] for the allowed multipoles. This beneficial feature allows sorting the magnets in the machine to carry out a local compensation of  $b_3$  [18].

#### VI. NOT ALLOWED MULTIPOLES

#### A. Even normal

Even normal multipoles are excited by a break down of the left-right symmetry in the coil. The LHC dipole features two-in-one collars that produce a systematic asymmetry between the left and the right part of each aperture. This asymmetry comes from both the mechanical structure of the collared coil, and from the magnetic structure due to asymmetry of the iron.

Here we aim at analyzing the coil asymmetries, and-therefore we consider the measurements at r.t. of the collared coil without iron yoke. We select data of Aperture 1 (on the right looking from the connection side), and similar results hold for Aperture 2. Data of Table III and Fig. 9 show that systematic

components are within one standard deviation, with the exception of  $b_2$  in Firm3. This means that the deformations of the two-in-one collars well preserve the left-right symmetry of each aperture. This is an expected result since the material chosen for the collar (austenitic steel) aims at minimizing the deformations of the structure. In Firm3, the negative systematic and the spread of  $b_2$  are driven by the imperfections of the collars, which are supplied by S2. This is proven (see [11], [12]) by the fact that contrary to Firm1 and 2, a strong correlation between apertures is observed Moreover, the expected average and spread of multipoles evaluated on the shasis of the collar dimensions (and on infinitely rigid collars) agree with the measured ones [11] Multipole  $b_4$  is well within targets and no trends are observed (see Fig. 10).

TABLE III: AVERAGE AND STANDARD DEVIATION OF EVEN NORMAL MULTIPOLES IN APERTURE I MEASURED IN THE COLLARED COILS AT R.T.

FOR FIRM 1 \$2

	b2		b4		b	6	b8	
Firm	μ	σ	μ	σ	μ	σ	μ	σ
1	-0.20	0.47	0.02	0.11	0.00	0.03	0.004	0.015
2	-0.22	0.44	-0.01	0.09	0.00	0.04	-0.003	0.019
3	-0.64	0.53	-0.06	0.09	0.00	0.03	-0.001	0.011
all	-0.35	0.53	-0.02	0.10	0.01	0.03	0.000	0.015

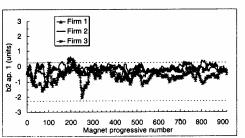


Fig. 9: Multipole b, measured at r.t. in the collared coil along the production: moving average for each dipole assembler (markers) and targets for systematic (dotted line).

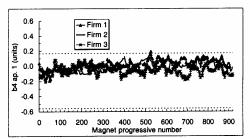


Fig. 10: Multipole b, measured at r.t. in the collared coil along the production: moving average for each dipole assembler (markers) and targets for systematic (dotted line).

The iron yoke is adding a  $b_2$  shift, of different sign in aperture of 2.5 units. This shift is very reproducible, with some systematic differences between the dipole assemblers, and aims at compensating the effect of cool down and saturation, thus optimizing  $b_2$  with respect to the beam dynamics targets.

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B. Oddskew FOR EXAMPLE,

Odd skew multipoles are excited by a break down of the symmetry associated to a rotation of  $\pi$  of the coil cross-section. Asystematic left-right asymmetry in the coil curing can excite these multipoles since the poles (i.e., the half coils) are assembled by rotating them by  $\pi$ . Iff the left part of the pole is longer than the right one, the coil obtains at itl of the coil midplane after assembly. Therefore, odd skew multipoles can give relevant information on the tolerances kept in the coil curing procedure. Another possible source of odd skew is a horizontal mismatch of the upper and the lower pole which could be due to collar imperfections [11], [12].

Average  $a_3$  has a non-zero systematic dipole assemblers of 0.3-0.4 units, with positive sign in Firm3 and negative in Firm1 and 2 (see Table IV and Fig. 11). A tilt of the mid-plane of 3 mrad gives 3.5 units of  $a_3$ , 0.3 units of  $a_5$ , and 0.07 units of  $a_7$  [19]. A comparison of these sensitivities with the  $a_3$  measurements of Table IV shows that the coil mid-plane tilt is kept within 0.15 mrad for the average of the production, within 0.4 mrad for the average of each dipole assembler, and within a spread (one stdev) of 0.3 mrad for the whole production. Indeed, it seems unlikely that the source of the systematic odd skews in Firm1 and in Firm2 is a mid-plane tilt, since  $\pm 0.3$  units of  $a_3$  would give rise to  $\pm 0.03$  units of  $a_5$ , which does not match the experimental value.

TABLE IV: AVERAGE AND STANDARD DEVIATION OF ODD SKEW MULTIPOLES IN

	EACH DIPOLE ASSEMBLER							BOTH APPETURES!			
		a3		<b>a</b> 5		a7		<b>a</b> 9			
I	Firm	μ	σ	μ	٩	μ	σ	μ	σ		
ĺ	1	-0.38	0.22	0.01	0.07	0.02	0.03	-0.023	0.018		
ı	2	-0.39	0.22	0.00	0.07	0.07	0.04	-0.003	0.025		
I	3	0.30	0.27	0.13	0.08	0.04	0.03	-0.009	0.014		
ľ	all	-0.16	0.42	0.05	0.09	0.04	0.04	-0.011	-0.011		

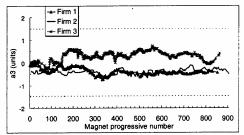


Fig. 11: Multipole a, measured at r.t. along the production: moving average for each dipole assembler (markers) and targets for systematic (dotted line)

Along the production, one observes that Firm3 has developed a positive systematic  $a_3$  since magnet  $150^{th}$  (see Fig. 11), and that a wide spike of negative  $a_3$  has been found in Firm2 between magnet  $400^{th}$  and  $550^{th}$  (see Fig. 12). In both cases, the sources of these trends are unknown.

The source of the systematic  $a_3$  in Firm3 could be a systematic left-right asymmetry of the collars, since-the used assembly procedure does not cancel this type of asymmetry (see [11] for details). This is confirmed by the fact that a good correlation is found between the two apertures of the same magnet, as expected from the assembly procedure. Moreover,

the  $a_3$  values deduced from the collar dimensions through a magneto-static model agree with the measured ones both for the average and for the sigma. The same conclusion cannot be drawn for  $a_3$  in Firm1 and 2, since the collar assembly procedure automatically cancels out the odd skews. Therefore the source of the systematic  $a_3$  is different  $a_3$ .

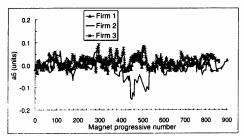


Fig. 12: Multipole a, measured at r.t. along the production: moving average for each dipole assembler (markers).

#### C. Even skew

Even skew multipoles are excited by a break down of the updown symmetry. A systematic up-down asymmetry can arise during the assembly of the collar around the coils (precollaring) in Firm2 and Firm3, where this operation is carried out horizontally, and during insertion of the locking rods under the press (collaring) in all assemblers. The possible source of asymmetry is the collar weight.

Average  $a_2$  has systematic constants close to zero in all

Average  $a_2$  has systematic competents close to zero in all dipole assemblers (within 0.2 units, i.e., ½ to 1/5 of the spread, see Table V and Fig. 13). Therefore, the vertical assembly procedure used in Firm1 does not improve the up-down symmetry with respect to the other two assemblers. A shift of the mid-plane of 0.1 mm gives 9.0 units of  $a_2$ , -1.0 units of  $a_4$ , and 0.27 units of  $a_6$  [19]. A comparison of these sensitivities with the  $a_2$  measurements gives the remarkable result that the coil mid-plane shift is kept within 1  $\mu$ m  $\mu$   $\mu$  for the average of the production, within 2  $\mu$ m for the average of each dipole assembler, and within a spread (one standard deviation) of 10  $\mu$ m for the whole production.

TABLE V: AVERAGE AND STANDARD DEVIATION OF EVEN SKEW MULTIPOLES

			IN EACH	ASSEMB	LEK	(PO ( II )		
	a2		a4		a6		a8	
Firm	μ	σ	μ.	σ	μ	σ	μ	σ
1	0.2	1.0	-0.03	0.22	0.00	0.06	0.000	0.017
2	-0.2	0.9	0.28	0.28	-0.01	0.08	0.025	0.024
3	0.2	0.7	-0.15	0.24	0.02	0.06	-0.010	0.019
all	0.0	0.9	0.03	0.30	0.00	0.07	0.005	0.025

On the other hand, a large non-zero systematic value of  $a_4$  is found in Firm2 (0.3 units), and a smaller one with opposite sign in Firm3 (see Fig. 14). The source of this systematic is not a coil mid-plane shift since 0.3 units of  $a_4$  would be associated to -2.7 units of  $a_2$ , which are not found. A more probable explanation for Firm2 is an inward radial movement of block 6 in the upper pole only of 50  $\mu$ m (see Fig. 15), which according to models gives rise to 0.3 units of  $a_2$  and to 0.27 units of  $a_4$ . This hypothesis is also supported by the fact that inward radial

movements of block 6, localized along the magnet longitudinal axis, in one quadrant only have been observed in Firm2 during the production (see [8] for more details).

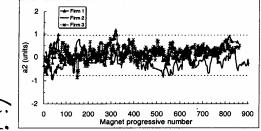


Fig. 13: Multipole a, measured at r.t. along the production: moving average for each dipole assembler (markers) and targets for systematic (dotted line)

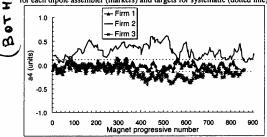


Fig. 14: Multipole a, measured at r.t. along the production: moving average for each dipole assembler (markers) and targets for systematic (dotted line).

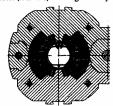


Fig. 15: Coil displacement of block 6 that could account for the systematic a measured at r.t. in Firm2 (not in scale)

Since these asymmetries are found already in the measurements of the collared coils, a simple way to cure them is to assemble half of the collared coils turned upside down in the cold masses. In this way the spread is higher but the systematic is zero. This solution has been successfully tested on one collared coil of Firm2 after one third of the production, when Firm3 had not yet developed a negative a4 and the situation was judged to be critical due to the absence of  $a_4$ correctors in the machine. It has not been implemented since the average  $a_4$  is within targets so far.

## VII. CONCLUSIONS

We discussed the mechanisms driving the field quality at room temperature in the main LHC dipoles, and its relation with tolerances and assembly procedures. For the bending strength we show that the spread is about 5 units, corresponding to a coil radius difference of 10 µm, and that the differences between assemblers are small due to an

accidental compensation of systematic differences in the transfer function and in the magnetic length.

The first order allowed component  $b_3$  is within targets and no differences between dipole assemblers are found. On the other hand, a signature of the dipole assembler is found in the higher order multipoles. This can be traced back to differences in the coil lay-out of 50 to 100 µm. Allowed components are not driven by the collar imperfections.

We showed that the even normal multipoles are within beam dynamic targets, and that the left-right symmetry of the assembly is well preserved, notwithstanding the two-in-one collars. For Firm3 there is strong evidence that  $b_2$  is driven by the collar imperfections. The analysis of skew multipoles shows that the up-down and rotational symmetry are well preserved. The coil mid-plane tilt is kept to zero within a fraction of mrad, and the shift is kept to zero within a few  $\mu\text{m}.$ A systematic component of  $a_4$  can be traced back to a nonsymmetric conductor displacement of 50 µm. Summarizing, the measured field quality at room temperature shows that the component and assembly tolerances set in the design have been successfully fulfilled in most cases. A few unexpected asymmetries or differences between assemblers are compatible with systematic displacements of less than 0.1 mm.

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